

# Search for CP and T violation in charm decays

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# Outline

- High impact physics
- CP violation in charm decays
- CP asymmetry: experimental status
- Dalitz plot analysis
- T-odd correlation
- Perspectives
- Conclusions

# High impact Physics

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Standard Model contains 2 key mysteries :

- the origin of the masses
- the existence of multiple fermion generations

The **first mystery** could be resolved by **LHC**

The **second mystery** appears to originate at higher mass scales **→** can only be studied **indirectly**

# High impact Physics

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CP violation, mixing and rare decays



may investigate the physics at these new scales !!

Why charm ?

Because in the charm sector the SM contributions to these effects are small



can provide **unique** information

# High impact Physics

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In addition **charm** is the unique probe of  
Up-type quark sector

but how **small** is small ?

- CP asymmetry  $\sim 10^{-3}$
- $D^0 - \bar{D}^0$  mixing  $\sim 10^{-7} \text{ -- } 10^{-10}$
- Rare decays  $\sim 10^{-9} \text{ -- } 10^{-19}$

High statistics instead of High Energy

Large window to search for new physics

# CP violation

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- **DIRECT CP VIOLATION:** refers to CP violation in meson decays where some CP violating phases necessarily appear in  $\Delta F = 1$  (decay) amplitudes
- **INDIRECT CP VIOLATION:** refers to CP violation in meson decays where the CP violating phases can all be chosen to appear in  $\Delta F = 2$  (mixing) amplitudes

# CP violation

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Three types of CP violation in Meson Decays:

- CP violation in mixing:

$$A_{SL} = \frac{\Gamma(D_{phys}^0(t) \rightarrow l^+ \nu X) - \Gamma(\bar{D}_{phys}^0(t) \rightarrow l^- \bar{\nu} X)}{\Gamma(D_{phys}^0(t) \rightarrow l^+ \nu X) + \Gamma(\bar{D}_{phys}^0(t) \rightarrow l^- \bar{\nu} X)}$$

- CP violation in decay:

$$A_{f^\pm} = \frac{\Gamma(D^+ \rightarrow f^+) - \Gamma(D^- \rightarrow f^-)}{\Gamma(D^+ \rightarrow f^+) + \Gamma(D^- \rightarrow f^-)}$$

- CP violation in the interference between decays with and without mixing:

$$A_{f_{CP}} = \frac{\Gamma(D_{phys}^0(t) \rightarrow f_{CP}) - \Gamma(\bar{D}_{phys}^0(t) \rightarrow f_{CP})}{\Gamma(D_{phys}^0(t) \rightarrow f_{CP}) + \Gamma(\bar{D}_{phys}^0(t) \rightarrow f_{CP})}$$

[Yosef Nir, Lecture at SLAC Summer Inst., hep-ph/9911321]

# CP violation

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- CP violating effects occur in a decay process only if the decay amplitude is the sum of 2 different parts, whose phases are made of a weak (CKM) and a strong (FSI) contribution:

$$\alpha = Ae^{i\delta_1} + Be^{i\delta_2}$$

The weak contributions to the phases change sign when going to the CP-conjugate process, while the strong ones do not.

The corresponding CP conjugate amplitude is:

$$\bar{\alpha} = A^*e^{i\delta_1} + B^*e^{i\delta_2}$$



# CP violation

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- Therefore the **CP** violating asymmetry in the decay rate will be :

$$A_{CP} = \frac{|\alpha|^2 - |\bar{\alpha}|^2}{|\alpha|^2 + |\bar{\alpha}|^2} = \frac{2\text{Im}(AB^*) \sin(\delta_2 - \delta_1)}{|A|^2 + |B|^2 + 2\text{Re}(AB^*) \cos(\delta_2 - \delta_1)}$$

Both factors in the numerator should be nonvanishing to have a **nonzero** effect. Moreover to have a sizeable asymmetry the moduli of the **2** amplitudes, A and B, should **not differ too much**

- In singly Cabibbo suppressed decays (SCSD) the penguin diagram contributions may provide a different phase for the second weak amplitude
- **FSI** provide a strong phase shift

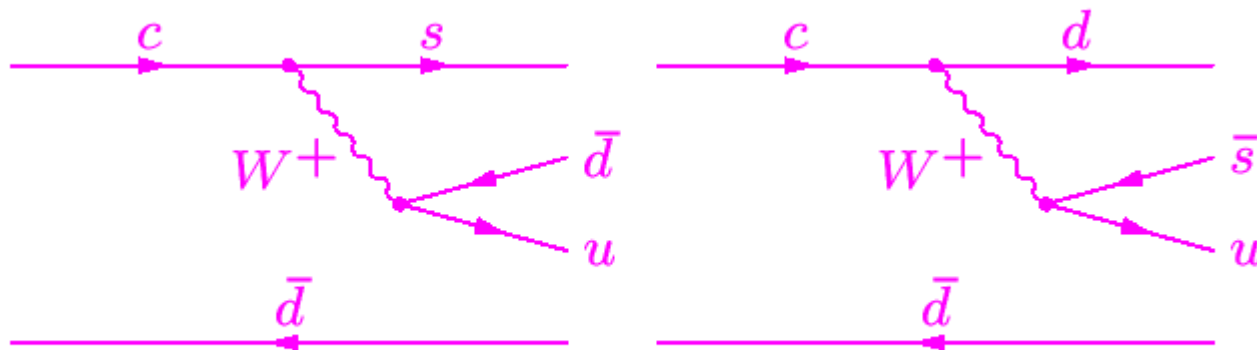
[Buccella et al., Phys.Lett.B302,319 (1993)]

# CP violation

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- For  $D^0$  decays there is the possibility of **indirect CP** violation due to **mixing**
- Buccella et al. predict **CP** violating decay asymmetries for charmed meson in the range:  
$$0.002 \% \longrightarrow 0.14 \%$$
- In the Standard Model **no** direct **CP** asymmetry can arise in Cabibbo allowed or DCS modes since they are driven by a single weak amplitude
- However in  $D \rightarrow K_s \pi$ 's **CP** asymmetries can arise in two different ways:
  - a) through the **CP** impurity in  $K_s$
  - b) through interference of two weak amplitudes

# CP violation



because one cannot differentiate between a  $K^0$  and a  $\bar{K}^0$  in the final state

- if New Physics intervenes through DCSD, then it would have the cleanest impact on  $D^+ \rightarrow K_{S,L} \pi^+$  ( Bigi and Sanda )

# CP violation

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“CP studies in charm transitions represent an almost zero background search for New Physics”

from “CP violation”(pag.252) by Bigi and Sanda

# CP violation search in experiments with large statistics

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- E791: fixed target experiment - hadroproduction of charm  
 $2.5 \times 10^5$  reconstructed charm
- CLEO: collider experiment -  $e^+ e^-$  at  $Y(4s)$   
 $9.0 \text{ fb}^{-1}$
- FOCUS: fixed target experiment - photoproduction of charm  
 $10^6$  reconstructed charm

# Experimental techniques

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In hadroproduction and photoproduction one needs to ratio to a Cabibbo allowed reference states in order to correct for **known production asymmetries**.

$$A_{CP} = \frac{\eta(D) - \eta(\bar{D})}{\eta(D) + \eta(\bar{D})}$$

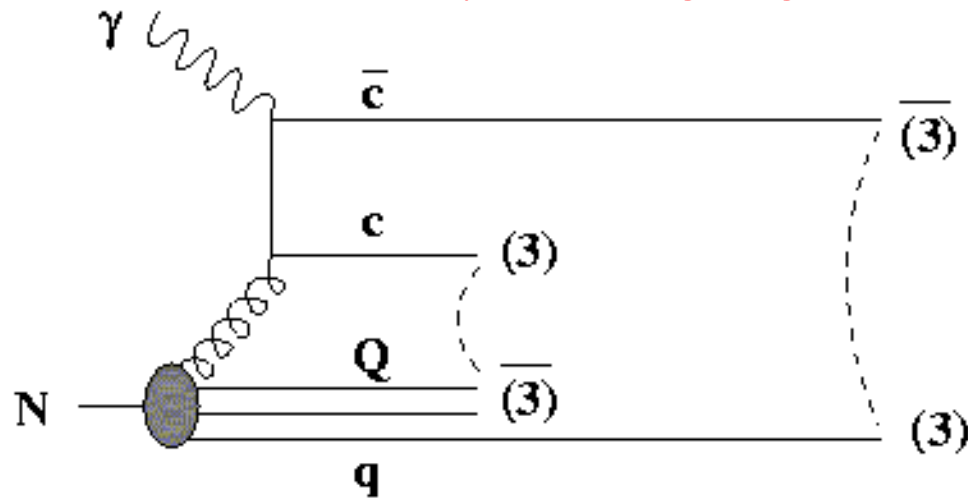
For example  $\eta$  for  $D^+ \rightarrow K^- K^+ \pi^+$  would be:  $\eta(D^+) = \frac{N(D^+ \rightarrow K^- K^+ \pi^+)}{N(D^+ \rightarrow K^- \pi^+ \pi^+)}$

For the  $D^0$ , the charm flavour is determined by tagging the charge of the bachelor **pion** from  $D^{*+} \rightarrow D^0 \pi^+$

# Experimental techniques

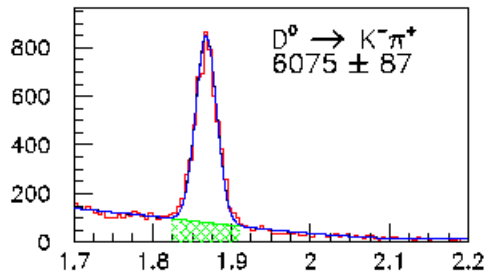
Why ratio the asymmetry at all?

Pythia string fragmentation model

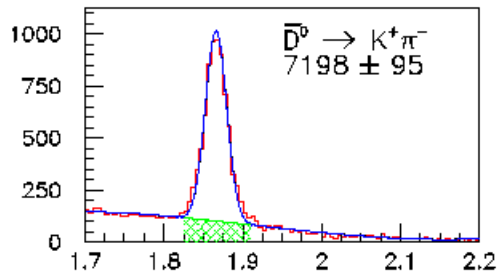


Need to correct for known production asymmetries  $\sim -5\%$  for photoproduced mesons. These likely arise from fragmentation dynamics. *eg* quark-diquark asymmetry

# E791 results

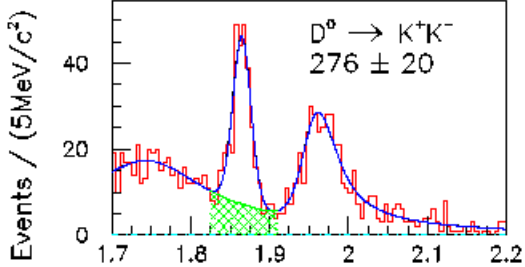


$M(K^-\pi^+)$  ( $\text{GeV}/c^2$ )

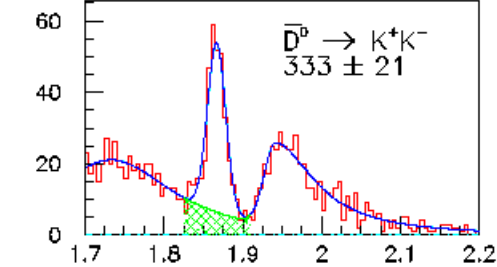


$M(K^+\pi^-)$  ( $\text{GeV}/c^2$ )

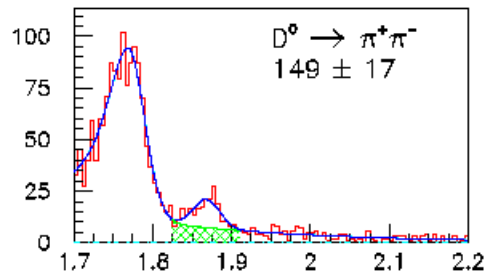
$$A_{CP} = \frac{\eta - \bar{\eta}}{\eta + \bar{\eta}}$$



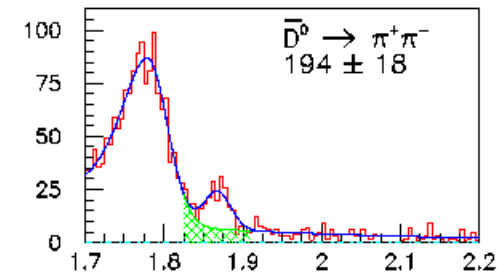
$M(K^+K^-)$  ( $\text{GeV}/c^2$ )



$M(K^+K^-)$  ( $\text{GeV}/c^2$ )



$M(\pi^+\pi^-)$  ( $\text{GeV}/c^2$ )



$M(\pi^+\pi^-)$  ( $\text{GeV}/c^2$ )

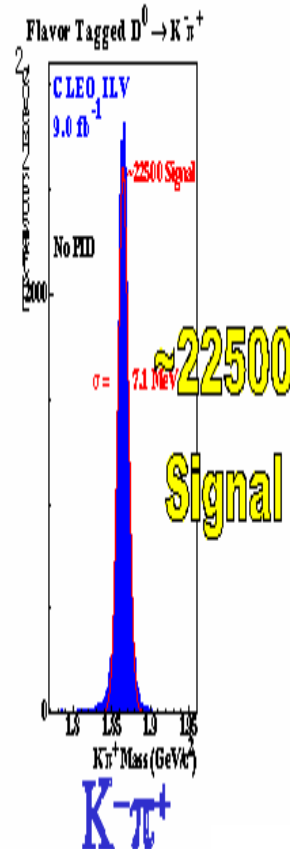
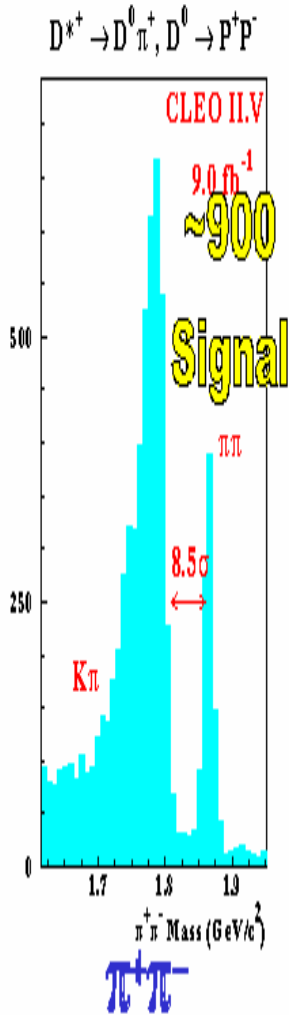
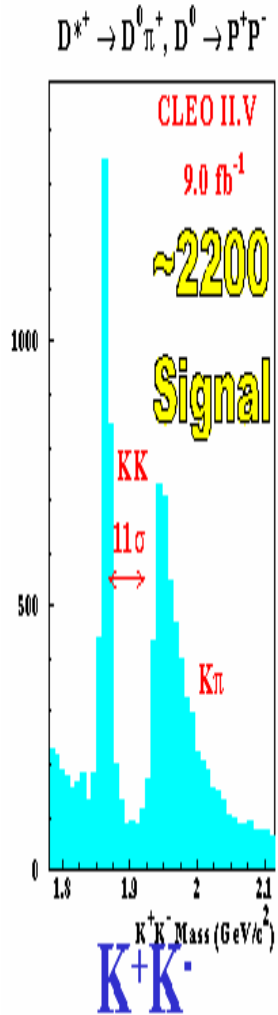
Decay mode	CP asymmetry (E791)
$D^0 \rightarrow K^- K^+$	$-0.010 \pm 0.049 \pm 0.012$
$D^0 \rightarrow \pi^- \pi^+$	$-0.049 \pm 0.078 \pm 0.030$
$D^+ \rightarrow K^- K^+ \pi^+$	$-0.014 \pm 0.029$
$D^+ \rightarrow \pi^- \pi^+ \pi^+$	$-0.017 \pm 0.042$

(PDG 2000)



# CLEO results

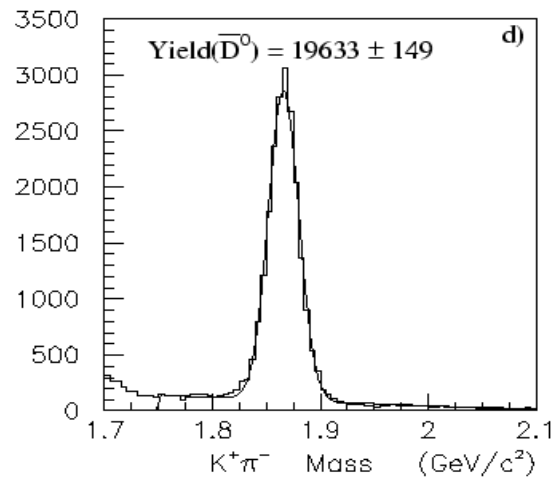
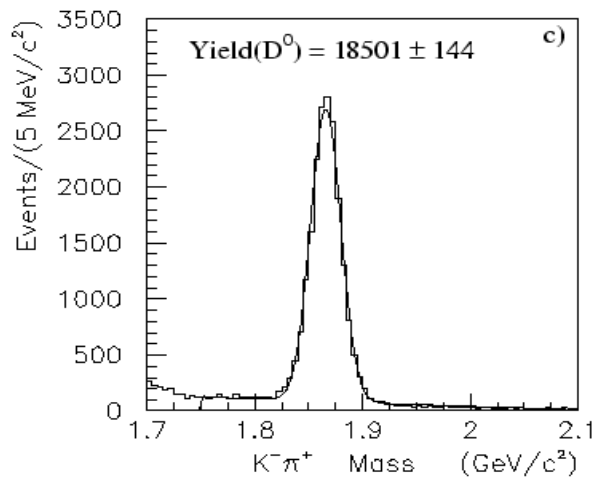
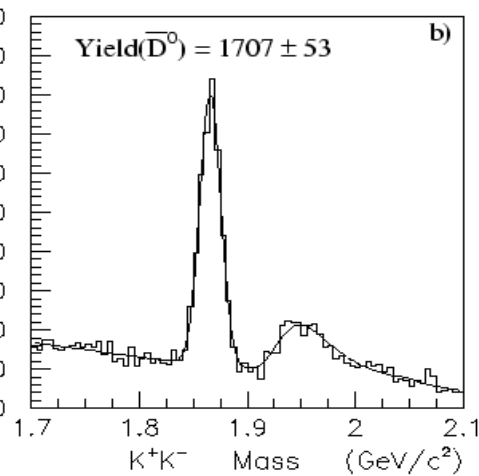
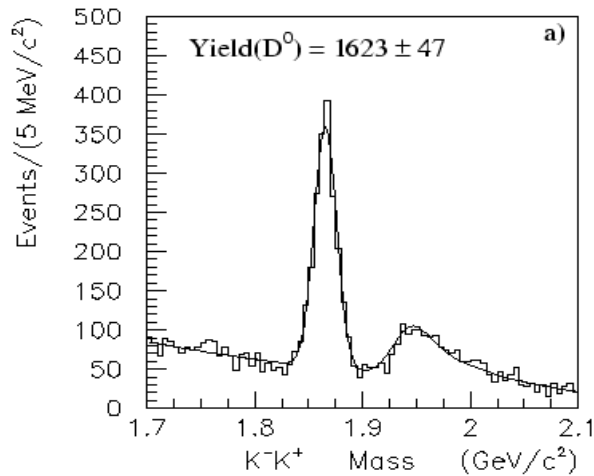
$$A_{CP} = \frac{N - \bar{N}}{N + \bar{N}}$$



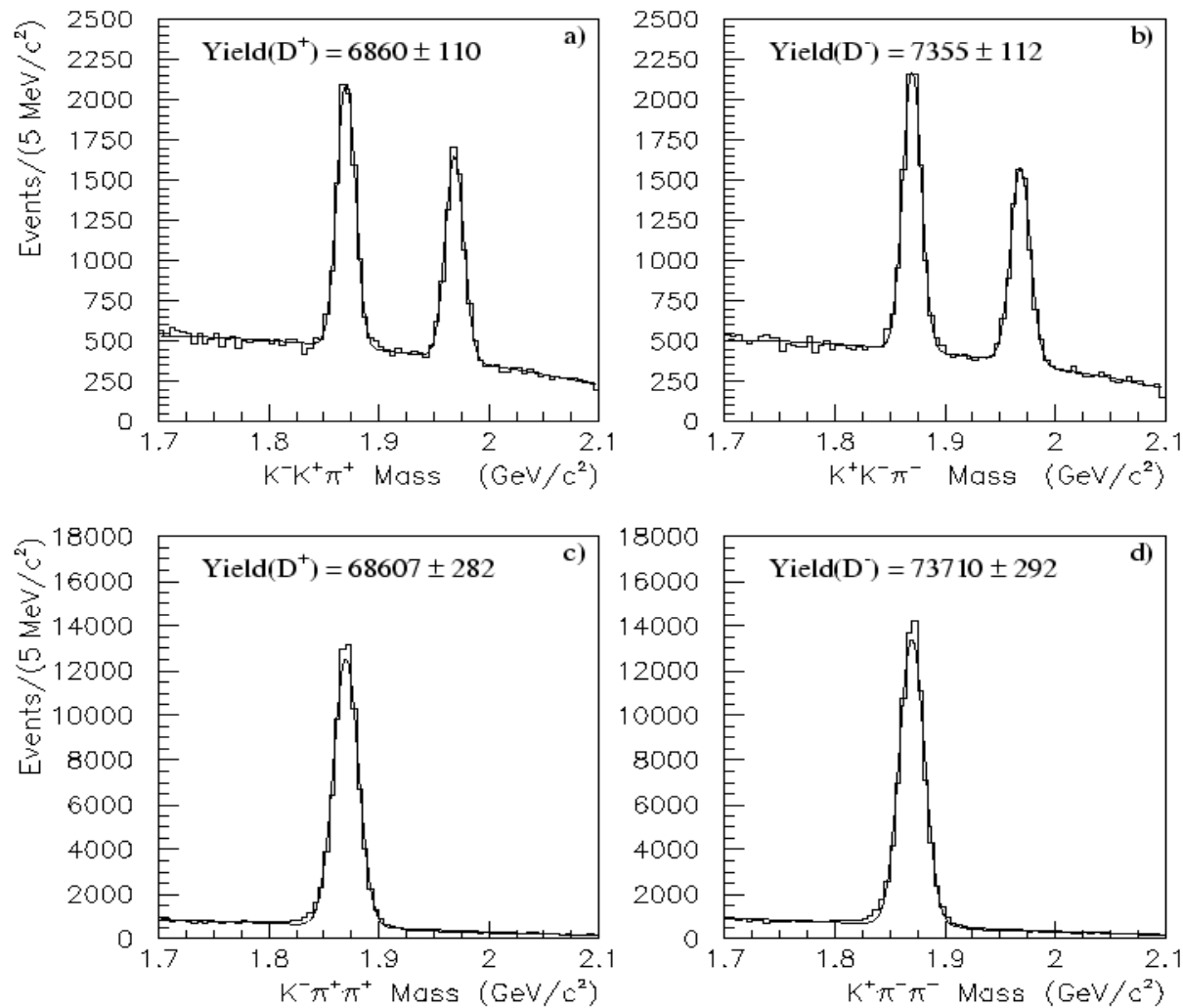
Decay mode	CP asymmetry (CLEO)
$D^0 \rightarrow K^- K^+$	$+0.0004 \pm 0.0218 \pm 0.0084$
$D^0 \rightarrow \pi^- \pi^+$	$+0.0194 \pm 0.0322 \pm 0.0084$
$D^0 \rightarrow K_S \pi^0$	$-0.018 \pm 0.030$

(Talk by Alex Smith at CIPANP 2000)

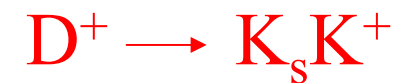
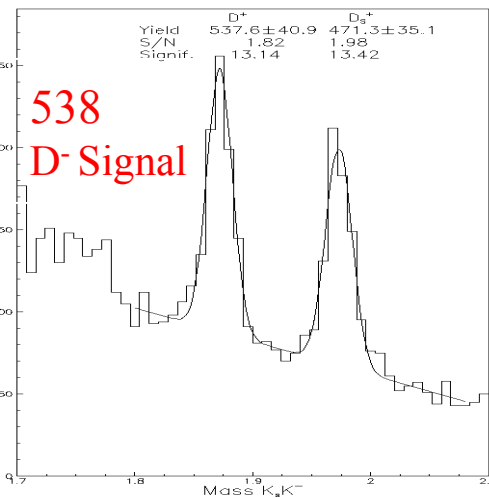
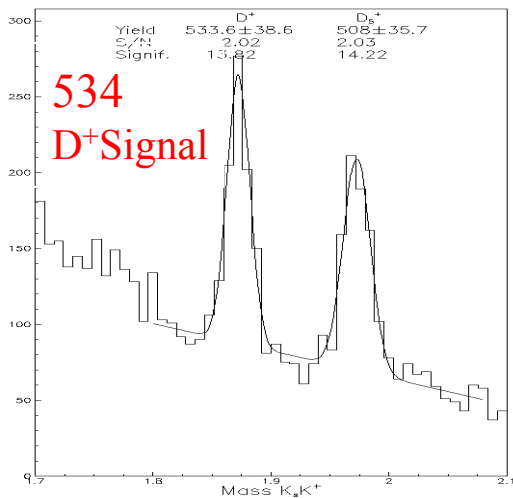
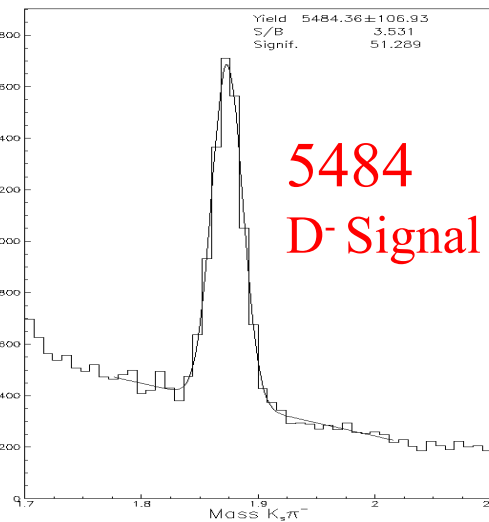
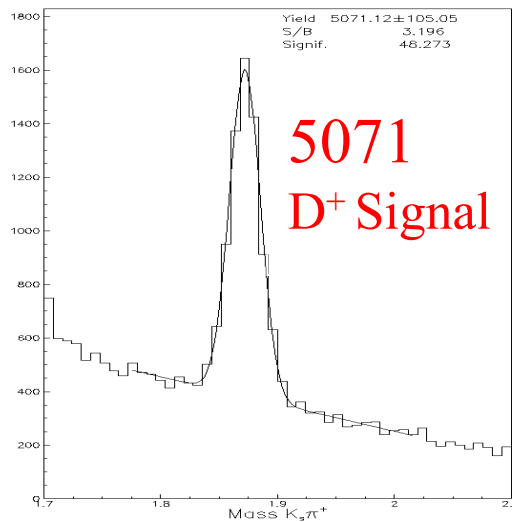
# CP violation search ( $D^0 \rightarrow K^- K^+$ )



# CP violation search ( $D^+ \rightarrow K^- K^+ \pi^+$ )



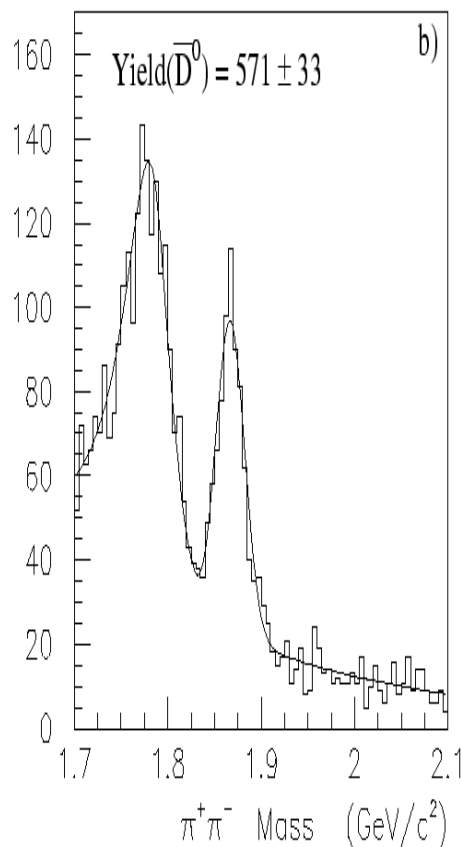
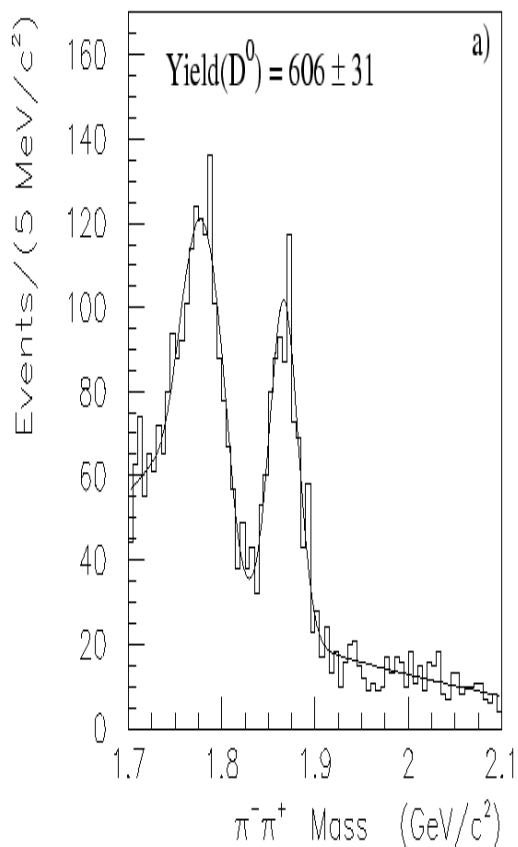
# CP violation search ( $D^+ \rightarrow K_s X^+$ )



(Talk by Brian O'Reilly  
at DPF 2000)



$$A_{CP} = \frac{\eta - \bar{\eta}}{\eta + \bar{\eta}}$$



Decay mode	CP asymmetry (FOCUS)
$D^0 \rightarrow K^- K^+$	$-0.001 \pm 0.022 \pm 0.015$
$D^0 \rightarrow \pi^- \pi^+$	$+0.048 \pm 0.039 \pm 0.025$
$D^+ \rightarrow K^- K^+ \pi^+$	$+0.006 \pm 0.011 \pm 0.005$
$D^+ \rightarrow K_S \pi^+$	$-0.016 \pm 0.015 \pm 0.009$
$D^+ \rightarrow K_S K^+$	$+0.069 \pm 0.060 \pm 0.015$

# Summary of **CP** asymmetry measurements

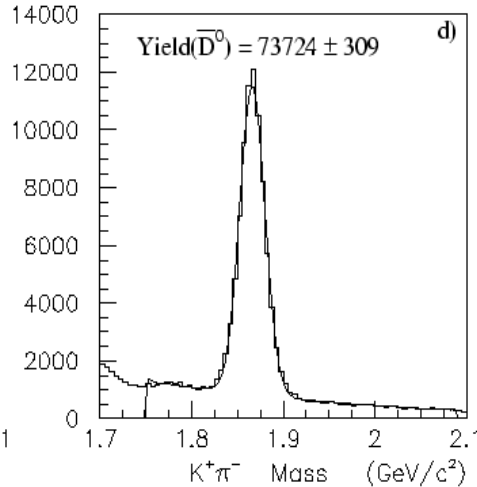
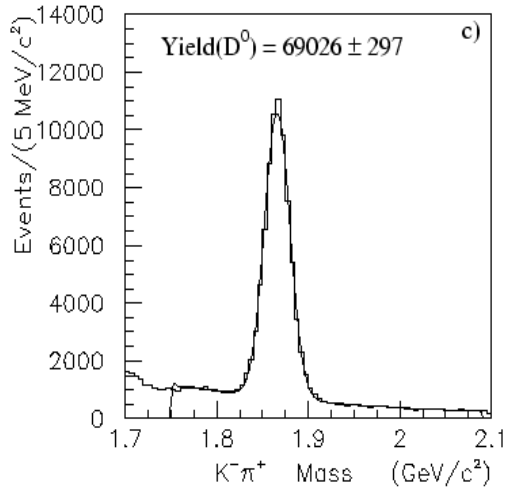
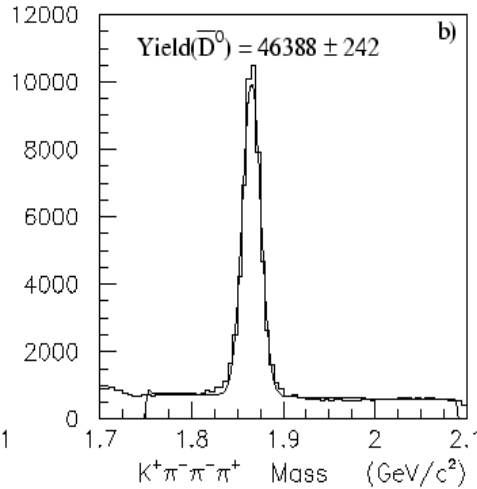
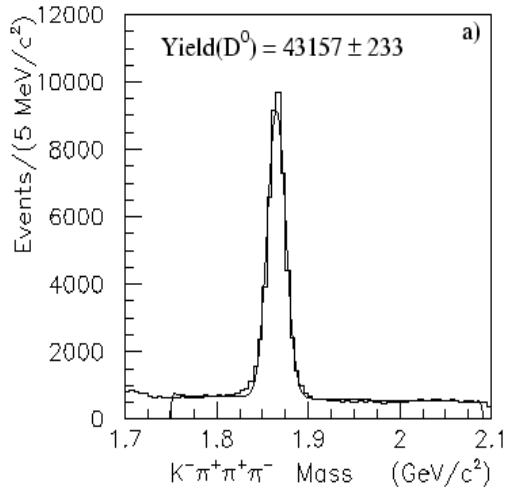
Decay mode	E791	CLEO	FOCUS
$D^0 \rightarrow K^- K^+$	$-0.010 \pm 0.049 \pm 0.012$	$+0.000 \pm 0.022 \pm 0.008$	$-0.001 \pm 0.022 \pm 0.015$
$D^0 \rightarrow \pi^- \pi^+$	$-0.049 \pm 0.078 \pm 0.030$	$+0.019 \pm 0.032 \pm 0.008$	$+0.048 \pm 0.039 \pm 0.025$
$D^0 \rightarrow K_S \pi^0$		$-0.018 \pm 0.030$	
$D^+ \rightarrow K^- K^+ \pi^+$	$-0.014 \pm 0.029$		$+0.006 \pm 0.011 \pm 0.005$
$D^+ \rightarrow \pi^- \pi^+ \pi^+$	$-0.017 \pm 0.042 \pm 0.005$		
$D^+ \rightarrow K_S \pi^+$			$-0.016 \pm 0.015 \pm 0.009$
$D^+ \rightarrow K_S K^+$			$+0.069 \pm 0.060 \pm 0.015$

- 1% level reached for some decay modes
- measured **CP** asymmetries are consistent with zero within errors
- **no evidence of CP violation**

# What's next.....

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- To search for New Physics beyond the Standard Model, one should look at the Cabibbo favored decay modes since they are **NOT** expected to exhibit **CP asymmetries**
- or Doubly Cabibbo Suppressed Decays....



$$A_{CP} = \frac{\eta - \bar{\eta}}{\eta + \bar{\eta}}$$

$$\eta = \frac{D^0 \rightarrow K^- \pi^+ \pi^- \pi^+}{D^0 \rightarrow K^- \pi^+}$$

$$A_{CP} = 0.0018 \text{ +/- } 0.0048(\text{stat.})$$

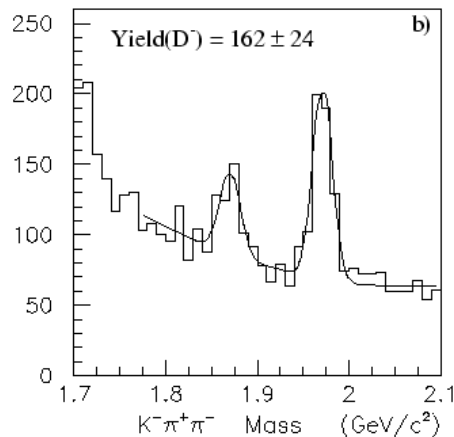
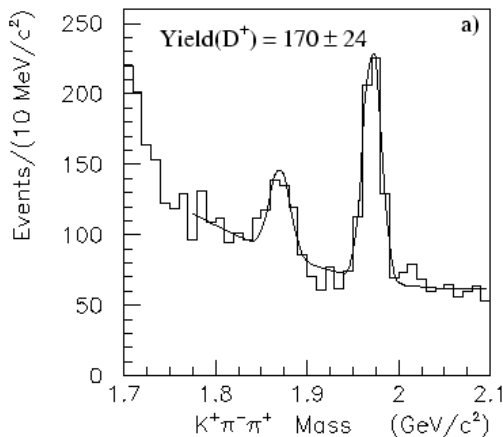
$$A_{CP} = 0.0032 \text{ +/- } 0.0048(\text{stat.})$$

(without MC efficiency correction)



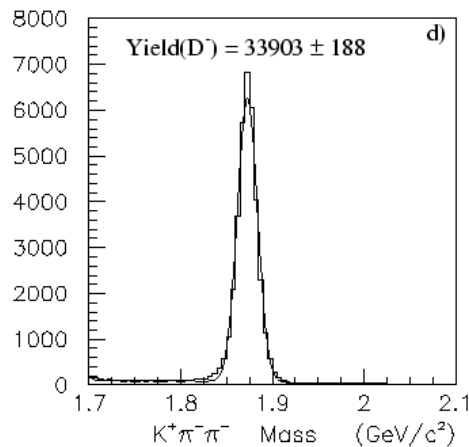
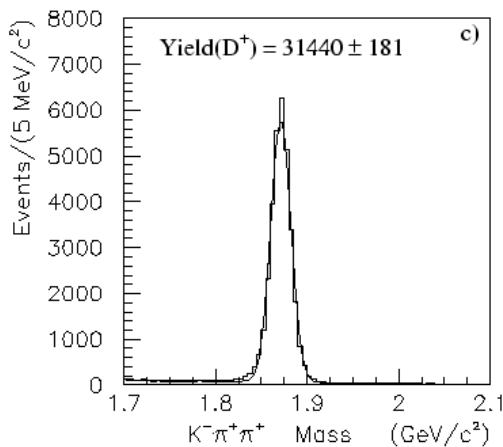
not even preliminary !!





$$A_{CP} = \frac{\eta - \bar{\eta}}{\eta + \bar{\eta}}$$

$$\eta = \frac{D^+ \rightarrow K^+ \pi^- \pi^+}{D^+ \rightarrow K^- \pi^+ \pi^+}$$



$$A_{CP} = 0.063 \pm 0.074 \text{ (stat.)}$$

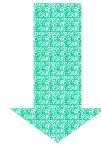


not even preliminary !!

# Dalitz plot analysis

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- Main advantage:  
complete information not only the branching ratio



DETERMINATION OF AMPLITUDE  
COEFFICIENTS AND PHASES

- Final state is the result of the interference of all the intermediate states

# Dalitz plot analysis

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- Decay amplitude: fit parameters  $a_i$  and  $\delta_i$

(3-body decays)

$$A(D) = \sum_i a_i e^{i\delta_i} B(abc | r_i)$$

where  $a, b, c$  label the final state particles and  $r$  the resonance

$$B(abc | r) = BW(a, b | r) S(a, c)$$

BW is the relativistic Breit - Wigner function :

$$BW(a, b | r) = \frac{F_D F_r}{M_r^2 - M_{ab}^2 - i\Gamma M_r}$$

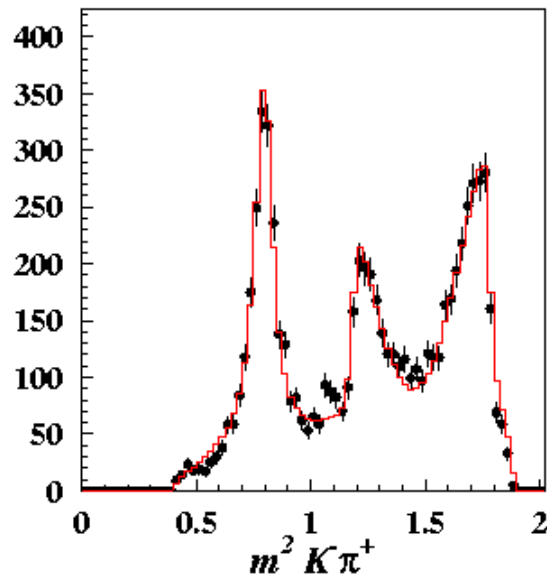
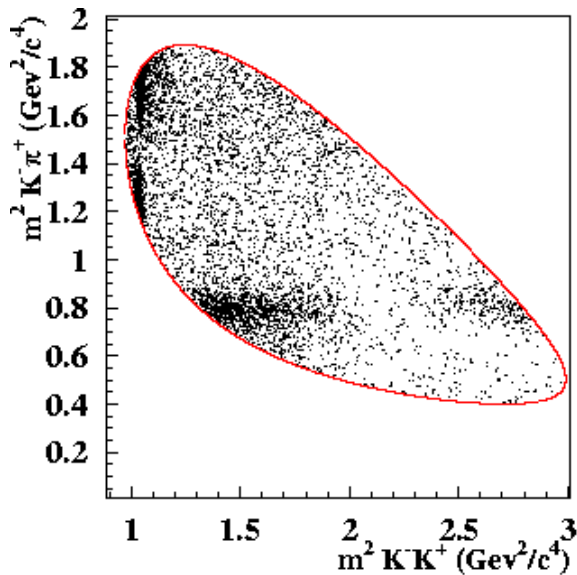
and  $S(a, c)$  is the spin term :      1    for spin 0 resonance

   -  $2\vec{c} \cdot \vec{a}$     for spin 1 resonance

.....

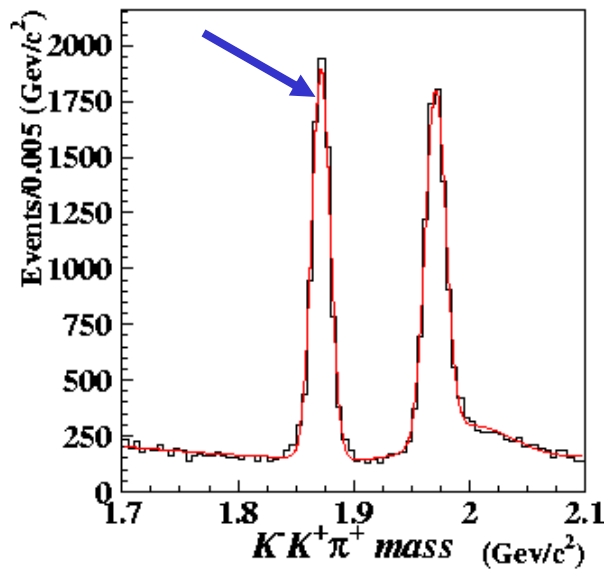
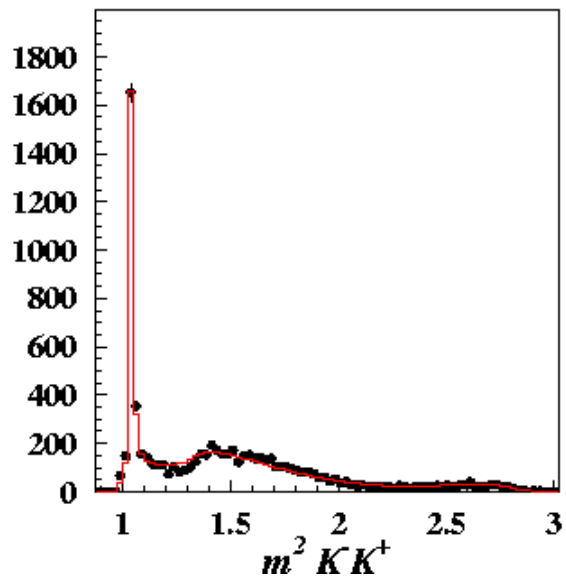


Preliminary results



Fit results

	Fit frac.	Phase (Deg)
$K^{*0}$ (892)	$0.31 \pm 0.01$	0 (fixed)
$K_0^*$ (1430)	$0.37 \pm 0.01$	$69 \pm 2$
$\phi$ (1020)	$0.29 \pm 0.01$	$-171 \pm 3$




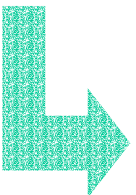
E687 published result

	Fit frac.	Phase (Deg)
$K^{*0}$ (892)	$0.30 \pm 0.02$	0 (fixed)
$K_0^*$ (1430)	$0.37 \pm 0.04$	$70 \pm 7$
$\phi$ (1020)	$0.29 \pm 0.03$	$-159 \pm 8$

(Talk by Sandra Malvezzi at CIPANP 2000)

# Dalitz plot analysis

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- $\delta_i = \sigma_i + \omega_i$   CP violating phase
-  CP conserving phase

under CP conjugation :

- $\bar{\delta}_i = \sigma_i - \omega_i$
- in general a difference between  $\delta_i$  and  $\bar{\delta}_i$  hints that CP is violated

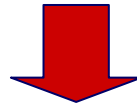
# Dalitz plot analysis

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Any difference in the reconstructed pattern, by Dalitz plot analysis, between  $D^+$  and  $D^-$  would be an evidence of **CP violation**

# Discrete symmetry P and T

Quantity	P	T
$r$	$-r$	$r$
$p$	$-p$	$-p$
$\sigma$	$\sigma$	$-\sigma$
$\sigma \cdot p$	$-\sigma \cdot p$	$\sigma \cdot p$
$\sigma \cdot (p_1 \times p_2)$	$\sigma \cdot (p_1 \times p_2)$	$-\sigma \cdot (p_1 \times p_2)$



*T-odd correlation*

$$C_T = v_1 \cdot (v_2 \times v_3)$$

where  $v_i$  is can be spin or momentum of a final state particle

# T-odd correlation

From I.I.Bigi 'Charm physics - Like Botticelli in the Sistine Chapel'  
arXiv:hep-ph/0107102 v1 (2001)

“ Consider, e.g.,  $D^0 \rightarrow K^- K^+ \pi^- \pi^+$ , where one can form a T-odd correlation with the momenta:

$$C_T = \langle p_{K^+} \circ (p_{\pi^+} \times p_{\pi^-}) \rangle$$

Under time reversal T one has  $C_T \rightarrow -C_T$  hence the name 'T-odd'.

Yet  $C_T \neq 0$  does not necessarily establishes T violation.

Since time reversal is implemented by an antiunitary operator,  $C_T \neq 0$  can be induced by FSI. While in contrast to the situation with partial width differences FSI are not required to produce an effect, they can act as an 'imposter' here, id est induce a T-odd correlation with T-invariant dynamics.

This ambiguity can unequivocally be resolved by measuring in  $D^0 \rightarrow K^- K^+ \pi^- \pi^+$ .

$$\bar{C}_T = \langle p_{K^-} \circ (p_{\pi^-} \times p_{\pi^+}) \rangle$$

Finding  $C_T \neq -\bar{C}_T$  establishes CP violation without further ado.”



# T-odd correlation

## Physical motivations

$$C_T = \mathbf{v}_1 \cdot (\mathbf{v}_2 \times \mathbf{v}_3)$$

(where  $v_i$  is a spin or a moment of a final particle)

Assuming  
CPT invariance

## T- odd correlation

( $C_T \rightarrow -C_T$  under time inversion)

**BUT**

To find  $C_T \neq 0$  doesn't imply CP violation

because

$C_T \neq 0$  can be produced through:

Weak phase  
(which violates CP)

Strong phase from F.S.I.  
(which doesn't violate CP)

**SOLUTION**

To study the CP conjugated process and to build  $C_T$  and  $(C_T)_{CP}$

If  $C_T \neq (C_T)_{CP} = -\bar{C}_T$  there is CP violation

# T-odd correlation

## Physical motivations

### How to compute CP asymmetries?

1. We build T-odd asymmetries using decay rates for a certain process and its CP conjugated process, as following:

$$A_T := \frac{\Gamma(C_T > 0) - \Gamma(C_T < 0)}{\Gamma(C_T > 0) + \Gamma(C_T < 0)}$$

where

$$A_T \doteq \sin(\delta_s + \phi_w)$$

$$(A_T)_{CP} := \frac{\Gamma((C_T)_{CP} > 0) - \Gamma((C_T)_{CP} < 0)}{\Gamma((C_T)_{CP} > 0) + \Gamma((C_T)_{CP} < 0)}$$

where

$$(A_T)_{CP} \doteq \sin(\delta_s - \phi_w)$$

NOT TRUE SIGNALS  
OF CP VIOLATION  
(because of F.S.I.)

2. We build a T-violation asymmetry, as following:

$$A_{Tviol} := \frac{1}{2} (A_T - (A_T)_{CP})$$

where

$$A_{Tviol} \doteq \cos(\delta_s) \sin(\phi_w)$$

TRUE SIGNAL  
OF CP VIOLATION  
(even in presence of F.S.I.)

To find  $A_{Tviol} \neq 0$  implies CP VIOLATION !

- W.Bensalem and D.London  
ArXiv:hep-ph/0005018 v1 (2000)

- G.Valencia Phys.Rev.D 39/11  
(1989) 3339

# T-odd correlation in the decay mode $D^0 \rightarrow K^- K^+ \pi^- \pi^+$

$D^0 \rightarrow K^- K^+ \pi^- \pi^+$  decay mode: final state with four different particles



We can build T-odd correlation:

$$C_T = \langle p_{K^+} \circ (p_{\pi^+} \times p_{\pi^-}) \rangle \quad \text{for the } D^0 \rightarrow K^- K^+ \pi^- \pi^+ \text{ decay mode}$$
$$\bar{C}_T = \langle p_{K^-} \circ (p_{\pi^-} \times p_{\pi^+}) \rangle \quad \text{for the } \bar{D}^0 \rightarrow K^- K^+ \pi^- \pi^+ \text{ decay mode}$$

$C_T \neq (C_T)_{CP} = -\bar{C}_T$  is an evidence of **CP VIOLATION**



We can build the T-odd asymmetries  $A_T$  and  $(A_T)_{CP}$   
and the T-violation asymmetry  $A_{Tviol}$

$A_{Tviol} \neq 0$  is an evidence of **CP VIOLATION**

# T-odd correlation in the decay mode $D^0 \rightarrow K^- K^+ \pi^- \pi^+$

- All the following results are **VERY PRELIMINARY**
- We are in touch with I. Bigi. He judged positively our analysis, even if he considers this work as a pilot study due to the limited statistics
- I. Bigi gave us some interesting advices and suggestions on how to proceed on this topic, looking for example at other decay modes:

$$D^0 \longrightarrow \bar{K}^0 \pi^0 \pi^+ \pi^-$$

$$D^+ \longrightarrow K^- K^+ \pi^0 \pi^+$$

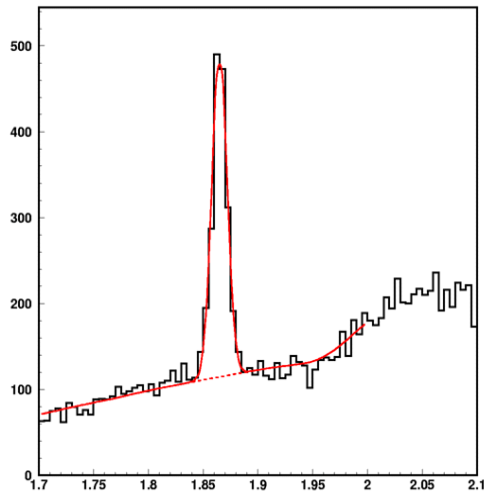
$$D^0 \longrightarrow K^- \pi^+ \pi^+ \pi^- \quad (\text{using the } \pi^+ \pi^- \text{ with largest invariant mass})$$

# $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ selection cuts

(following Alberto Reis analysis)

## Old cuts:

- $L/S > 10$
- $Mulprim > 1$
- $Confidence\ Level > 0.10$
- $Iso2 < 0.005$
- $Pionicity > 0.5$
- $Kaonicity > 2$

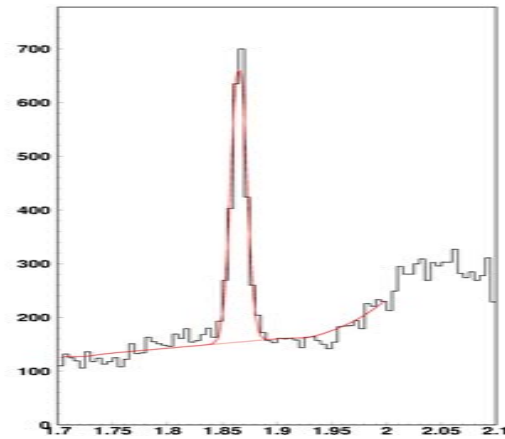


$Y = 1304 \pm 53$   
 $S/N = 3.2 \pm 0.3$

## New cuts:

- $L/S > 10$
- $Mulprim > 1$
- $Confidence\ Level > 0.01$
- $Iso2 < 0.10$
- $\pi^-$  consistency  $< 5$
- $Kaonicity > 2$
- **Score**  $> 0.15$
- $Zout > 0$

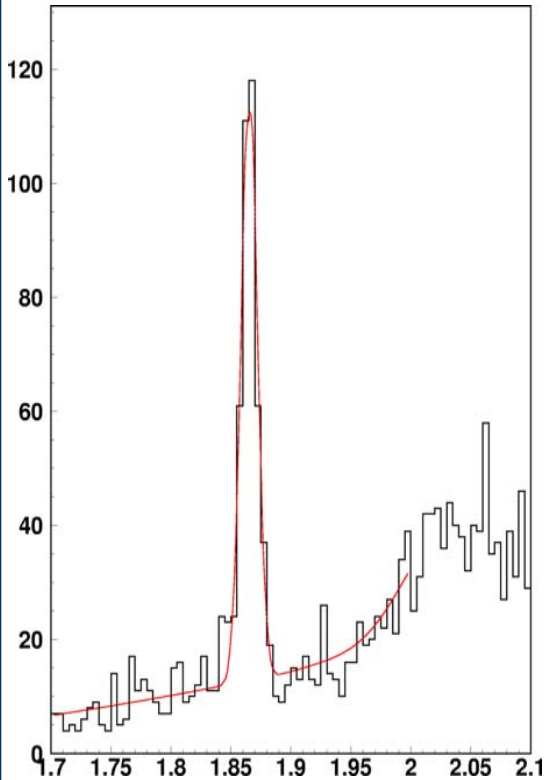
$$score = \frac{1}{3} \left( \frac{L/\sigma}{(L/\sigma)_{max}} + \frac{C.L.}{(C.L.)_{max}} + \frac{Zout}{(Zout)_{max}} \right)$$



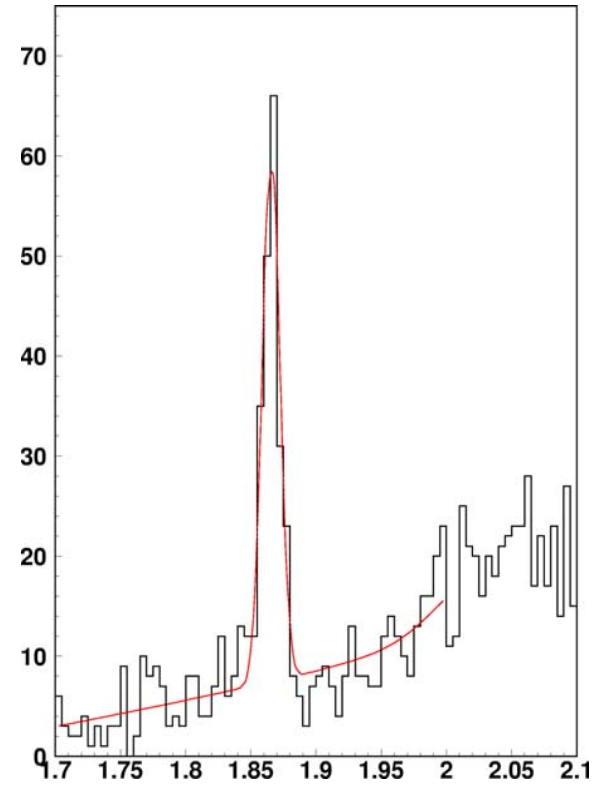
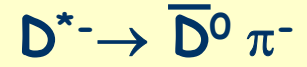
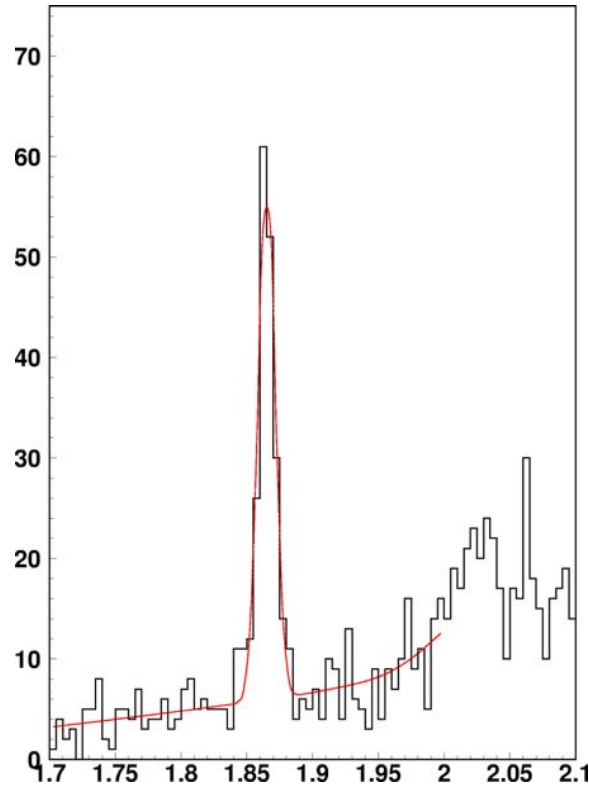
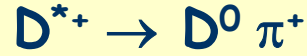
$Y = 1831 \pm 59$   
 $S/N = 2.9 \pm 0.4$

# D\*-Tag

D\*-Tag sample



We can distinguish the particle from the antiparticle using the D\*-Tag:

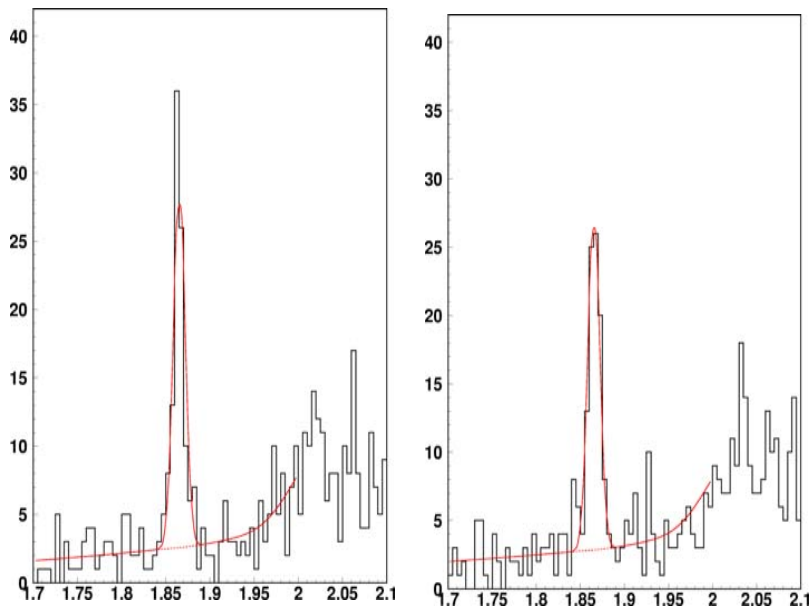


# T-odd correlation in the decay mode $D^0 \rightarrow K^- K^+ \pi^- \pi^+$

$$D^0 \rightarrow K^- K^+ \pi^- \pi^+$$

$$C_T = \langle p_{K^+} \circ (p_{\pi^+} \times p_{\pi^-}) \rangle$$

T-odd correlation



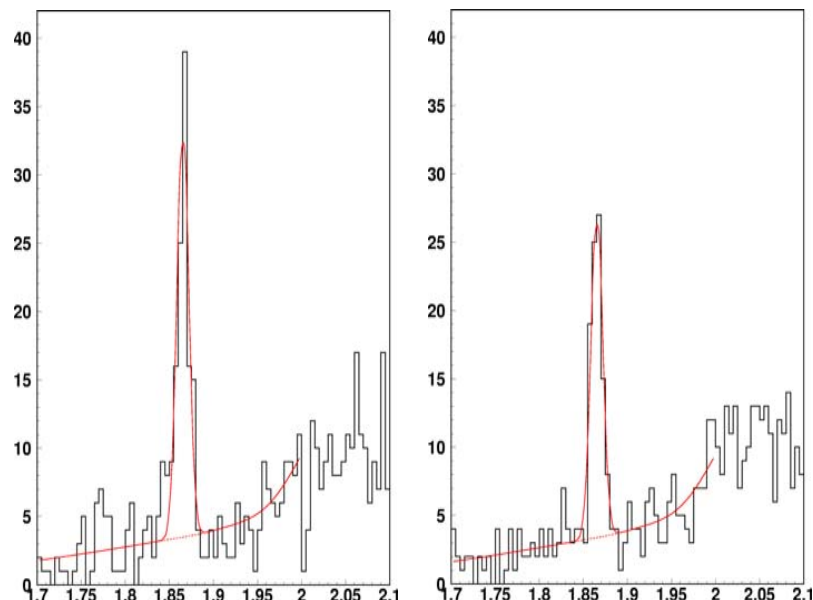
$$C_T > 0$$

$$C_T < 0$$

$$\bar{D}^0 \rightarrow K^- K^+ \pi^- \pi^+$$

$$\bar{C}_T = \langle p_{K^-} \circ (p_{\pi^-} \times p_{\pi^+}) \rangle$$

T-odd correlation



$$\bar{C}_T > 0$$

$$\bar{C}_T < 0$$

# T-odd correlation in the decay mode $D^0 \rightarrow K^- K^+ \pi^- \pi^+$

decay mode	request	yield
$D^0 \rightarrow K^- K^+ \pi^- \pi^+$	$C_T > 0$	$88 \pm 10$
	$C_T < 0$	$82 \pm 10$
$\bar{D}^0 \rightarrow K^- K^+ \pi^- \pi^+$	$\bar{C}_T > 0$	$101 \pm 11$
	$\bar{C}_T < 0$	$80 \pm 10$



$$A_{T\text{viol}} = 0.075 \pm 0.064$$



**NO EVIDENCE of CP VIOLATION**



# Perspectives

In the end, a very promising future:

- new results are coming from **FOCUS** and **CLEO**
- and recently from **BaBar**, **BELLE**
- asymptotically from **BTeV** and **LHC-B**

# Perspectives

Exp.	Beam	Lum.	Int. L (10**7 sec)	Cross Sec. cc	Cross Sec. bb	#evts cc prod.	#evts bb prod.	# evt cc ric. <b>estim.</b>	Physics start date
BABAR	e+e- Y(4s) asym	3 x 10**33 cm-2s-1	30 fb-1		1.2 nb		3.6 x 10**7	3.6 x 10**6	1999
BELLE	e+e- Y(4s) asym.	3 x 10**33 cm-2 s-1	30 fb-1		1.2 nb		3.6 x 10**7	3.6 x 10**6	1999
COMPASS	FT $\pi$ -Cu	~ 1 x 10**32 cm-2 s-1	1 fb-1	~ 10 $\mu$ b		10**9		~ 5 x 10**6	2001
CESR/ CLEO-C (?)	e+e- 3-5 Gev	2 x 10**32 cm-2 s-1	2 fb-1	~ 10 nb		2 x 10**7		2 x 10**6	2003
BTeV	ppbar 2TeV	2 x 10**32 cm-2 s-1	2 fb-1	> 500 $\mu$ b	100 $\mu$ b	10**12	2 x 10*11	10**9	>2005
LHC-b	pp 14 TeV	2 x 10**32 cm-2 s-1	2 fb-1	> 500 $\mu$ b	500 $\mu$ b	> 10**12	10**12		>2005

# Conclusions

- No **CP** violation has been reported in the charm sector
- However the role of charm remains **unique** in probing SM and looking for New Physics
- the **c** quark is the only “**u-type**” quark for which the decay modes can be studied
- complementary to the “**d-type**” sector investigation

# Conclusions

- Advantage of using Dalitz plot analysis



direct access to the **phases**

- **Possibility of studying T-odd correlation, in some particular decay modes, in order to determine T-violation**

# Two bodies Branching Ratio evaluation

This result:

$$\frac{\Gamma(D^0 \rightarrow K^+ K^-)}{\Gamma(D^0 \rightarrow K^- \pi^+)} = 0.1015 \pm 0.0013 \pm 0.0010$$

$$\frac{\Gamma(D^0 \rightarrow \pi^+ \pi^-)}{\Gamma(D^0 \rightarrow K^- \pi^+)} = 0.0374 \pm 0.0009 \pm 0.0004$$

$$\frac{\Gamma(D^0 \rightarrow K^+ K^-)}{\Gamma(D^0 \rightarrow \pi^- \pi^+)} = 2.71 \pm 0.09$$

Previous results:

0.1109±0.0033 (PDG)  
0.109±0.003±0.003 (E791)  
0.116±0.007±0.007 (CLEO)  
0.109±0.007±0.009 (E687)

0.0397±0.0021 (PDG)  
0.040±0.002±0.003 (E791)  
0.043±0.007±0.003 (E687)  
0.0348±0.0030±0.0023 (CLEO)

2.75 ± 0.15 ± 0.16 (E791)  
2.53 ± 0.46 ± 0.19 (E687)

# Four bodies Branching Ratio evaluation

This result:

$$\frac{\Gamma(D^0 \rightarrow K^+ K^- \pi^- \pi^+)}{\Gamma(D^0 \rightarrow K^- \pi^+ \pi^- \pi^+)} = 0.0297 \pm 0.0010 \pm 0.0008$$

$$\frac{\Gamma(D^0 \rightarrow \pi^+ \pi^- \pi^+ \pi^-)}{\Gamma(D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-)} = 0.0866 \pm 0.0012 \pm 0.0005$$

$$\frac{\Gamma(D^0 \rightarrow K^+ K^- \pi^- \pi^+)}{\Gamma(D^0 \rightarrow \pi^- \pi^+ \pi^- \pi^+)} = 0.34 \pm 0.01$$

Previous results:

$$0.0334 \pm 0.0028 \text{ (PDG)}$$
$$0.0313 \pm 0.0037 \pm 0.0036 \text{ (E791)}$$
$$0.035 \pm 0.004 \pm 0.002 \text{ (E687)}$$

$$0.098 \pm 0.006 \text{ (PDG)}$$
$$0.095 \pm 0.007 \pm 0.002 \text{ (E687)}$$

$$0.34 \pm 0.04 \text{ (PDG)}$$

# Sum of the related 2 and 4 bodies final states

(Theoretical suggestion by I.Bigi)

$$\frac{\Gamma(D^0 \rightarrow K^- K^+ + D^0 \rightarrow K^- K^+ \pi^- \pi^+)}{\Gamma(D^0 \rightarrow \pi^- \pi^+ + D^0 \rightarrow \pi^- \pi^+ \pi^- \pi^+)} = 0.80 \pm 0.04$$



PDG result  
 $0.8 \pm 0.1$

The ratio is close to 1



Adding related decay modes in 2 and 4 bodies gives

**REDUCTION OF THE  $SU(3)_{fl}$  SIMMETRY BREAKING**

Actually a complete analysis of the 4-body decay rates requires the study of the subresonant structure.

# Isospin analysis of $D \rightarrow \pi \pi$ decay channel

## Theoretical study

$$D^0 \rightarrow \pi^- \pi^+$$

$$D^0 \rightarrow \pi^0 \pi^0$$

$$D^+ \rightarrow \pi^+ \pi^0$$

Decay modes

Isospin  
space



$$\left| \frac{1}{2}, -\frac{1}{2} \right\rangle \rightarrow \sqrt{\frac{1}{3}} |2,0\rangle + \sqrt{\frac{2}{3}} |0,0\rangle$$

$$\left| \frac{1}{2}, -\frac{1}{2} \right\rangle \rightarrow \sqrt{\frac{2}{3}} |2,0\rangle - \sqrt{\frac{1}{3}} |0,0\rangle$$

$$\left| \frac{1}{2}, +\frac{1}{2} \right\rangle \rightarrow |2,+1\rangle$$

Decay modes in Isospin space

Hamiltonian

$$H = A_{3/2} \left| \frac{3}{2}, +\frac{1}{2} \right\rangle + A_{1/2} \left| \frac{1}{2}, +\frac{1}{2} \right\rangle$$

Lipkin, Nir, Quinn, Snyder  
*Phys.Rev.D* 44:1454 (1991)

$$A^{+-} = (\sqrt{2} A_0 + A_2) / \sqrt{3}$$

$$A^{00} = (A_0 + \sqrt{2} A_2) / \sqrt{3}$$

$$A^{+0} = A_2 (\sqrt{3} / \sqrt{2})$$

Decay amplitudes

Isospin  
space



where

$$\begin{aligned} A_2 &:= \frac{1}{\sqrt{2}} A_{3/2,2} \\ A_0 &:= \frac{1}{\sqrt{2}} A_{1/2,0} \end{aligned}$$

$A_{if}$  Isospin  
amplitudes

$A_{it,If}$  Transition  
amplitudes

$$\begin{aligned} \left| \frac{A_2}{A_0} \right|^2 &= \frac{\frac{2}{3} |A^{+0}|^2}{|A^{+-}|^2 - \frac{2}{3} |A^{+0}|^2} \\ \cos(\delta_2 - \delta_0) &= \frac{3|A^{+-}|^2 - 6|A^{00}|^2 + 2|A^{+0}|^2}{4\sqrt{2}|A^{+0}|^2 \sqrt{\frac{3}{2}(|A^{00}|^2 + |A^{+-}|^2) - |A^{+0}|^2}} \end{aligned}$$

Isospin amplitudes and phases

Cleo collaboration  
*Phys.Rev. Lett.* 71 (13)/1973 (1993)



# Isospin analysis of $D \rightarrow \pi \pi$ decay channel

## Experimental results

$$\Gamma^{+-} = (0.227 \pm 0.009) 10^{-14} \text{ Gev}$$

$$\Gamma^{00} = (0.156 \pm 0.044) 10^{-14} \text{ Gev}$$

$$\Gamma^{+0} = (0.133 \pm 0.035) 10^{-14} \text{ Gev}$$

Where  $\rightarrow$

$$\Gamma^{+-} = \frac{\hbar}{\tau_{D^0}} \left( \frac{\Gamma(K^- \pi^+)}{\Gamma_{tot}} \right) \left( \frac{\Gamma(\pi^- \pi^+)}{\Gamma(K^- \pi^+)} \right)$$

$$\Gamma^{00} = \frac{\hbar}{\tau_{D^0}} \left( \frac{\Gamma(\pi^0 \pi^0)}{\Gamma_{tot}} \right)$$

$$\Gamma^{+0} = \frac{\hbar}{\tau_{D^0}} \left( \frac{\Gamma(\pi^+ \pi^0)}{\Gamma_{tot}} \right)$$

Our study  $\rightarrow$  PDG

$$|A_2/A_0| = 0.64 \pm 0.13$$

$$\cos(\Delta\delta) = 0.14 \pm 0.17$$

$$\Delta\delta = (81.9 \pm 9.8)^\circ$$

This result

$\longleftrightarrow$

$$|A_2/A_0| = 0.72 \pm 0.13 \pm 0.11$$

$$\cos(\Delta\delta) = 0.14 \pm 0.13 \pm 0.09$$

Cleo result

### Conclusions:

- Isospin amplitudes are of the same order
- Large phase difference between the two isospin amplitudes



Evidence of elastic F.S.I.

# Isospin analysis of $D \rightarrow K K$ decay channel

## Theoretical study

$$D^0 \rightarrow K^- K^+$$

$$D^0 \rightarrow K^0 \bar{K}^0$$

$$D^+ \rightarrow K^+ K^0$$

Decay modes

Isospin  
space



$$\begin{aligned} \left| \frac{1}{2}, -\frac{1}{2} \right\rangle &\rightarrow \sqrt{\frac{1}{2}} (|1,0\rangle + |0,0\rangle) \\ \left| \frac{1}{2}, -\frac{1}{2} \right\rangle &\rightarrow \sqrt{\frac{1}{2}} (|1,0\rangle - |0,0\rangle) \\ \left| \frac{1}{2}, +\frac{1}{2} \right\rangle &\rightarrow |1,+1\rangle \end{aligned}$$

Decay modes in Isospin space

$$H = A_{1/2} \left| \frac{1}{2}, +\frac{1}{2} \right\rangle$$

Hamiltonian

$$A^{+-} = (A_1 + A_0) / \sqrt{2}$$

$$A^{00} = (A_1 - A_0) / \sqrt{2}$$

$$A^{+0} = A_1 \sqrt{2}$$

Decay amplitudes

$A_{If}$  Isospin  
amplitudes

Isospin  
space



where

$$\begin{aligned} A_1 &:= \frac{1}{\sqrt{2}} A_{1/2,1} \\ A_0 &:= \frac{1}{\sqrt{2}} A_{1/2,0} \end{aligned}$$

$$\left| \frac{A_1}{A_0} \right|^2 = \frac{|A^{+0}|^2}{2|A^{+-}|^2 + 2|A^{00}|^2 - |A^{+0}|^2}$$

$$\cos(\delta_1 - \delta_0) = \frac{|A^{+-}|^2 - |A^{00}|^2}{\sqrt{|A^{+0}|^2} \sqrt{2|A^{+-}|^2 + 2|A^{00}|^2 - |A^{+0}|^2}}$$

Isospin amplitudes and phases

$A_{it,If}$  Transition  
amplitudes

# Isospin analysis of $D \rightarrow K K$ decay channel

## Experimental results

$$\Gamma^{+-} = (0.616 \pm 0.019) 10^{-14} \text{ GeV}$$

$$\Gamma^{00} = (0.461 \pm 0.063) 10^{-14} \text{ GeV}$$

$$\Gamma^{+0} = (0.103 \pm 0.029) 10^{-14} \text{ GeV}$$

Where  $\rightarrow$

$$\Gamma^{+-} = \frac{\hbar}{\tau_{D^0}} \left( \frac{\Gamma(K^- \pi^+)}{\Gamma_{tot}} \right) \left( \frac{\Gamma(K^- K^+)}{\Gamma(K^- \pi^+)} \right)$$

$$\Gamma^{00} = \frac{\hbar}{\tau_{D^0}} \left( \frac{\Gamma(K^0 \bar{K}^0)}{\Gamma_{tot}} \right)$$

$$\Gamma^{+0} = \frac{\hbar}{\tau_{D^0}} \left( \frac{\Gamma(K^+ K^0)}{\Gamma_{tot}} \right)$$

Our study  
PDG

$$|A_1/A_0| = 0.69 \pm 0.07$$

$$\cos(\Delta\delta) = 0.76 \pm 0.07$$

$$\Delta\delta = (40.3 \pm 6.3)^\circ$$

This result



$$|A_1/A_0| = 0.61_{-0.10}^{+0.11}$$

$$\cos(\Delta\delta) = 0.88_{-0.08}^{+0.10}$$

Cleo result

### Conclusions:

- Isospin amplitudes are of the same order
- Phase difference between the two isospin amplitudes (but smaller than the previous one)



Evidence of elastic F.S.I.

# Effects of elastic F.S.I. on B.R. ( $D^0 \rightarrow K^- K^+$ ) / ( $D^0 \rightarrow \pi^- \pi^+$ )

Elastic F.S.I. - rotation in Isospin space



$$\frac{\Gamma(D^0 \rightarrow K^- K^+)}{\Gamma(D^0 \rightarrow \pi^- \pi^+)} = 2.71 \pm 0.09$$



$$\frac{\Gamma(D^0 \rightarrow K^+ K^-) + \Gamma(D^0 \rightarrow K^0 \bar{K}^0)}{\Gamma(D^0 \rightarrow \pi^- \pi^+) + \Gamma(D^0 \rightarrow \pi^0 \pi^0)} = 1.97 \pm 0.20$$

Not affected by elastic F.S.I.

Large SU(3) breaking

Large (but inferior) SU(3) breaking

## Conclusions:

- Summing over Isospin rotated channels, SU(3) breaking is reduced.
- The effect of elastic F.S.I. on B.R. ( $K^- K^+$ ) / ( $\pi^- \pi^+$ ) is significant, but is unable to explain the discrepancy between experimental results and theoretical predictions ( $\Gamma(K^- K^+) / \Gamma(\pi^- \pi^+) < 1.4$ )



Possible solution : inelastic F.S.I.

# References for Isospin analysis

M.Gronau, D.London

*Phys Rev. Lett.* 65,, 27:3381 (1990)

H.J.Lipkin, Y.Nir, H.R.Quinn, A.E.Syder D.London

*Phys. Rev. D* 44,5:1454 (1991)

Cleo Collaboration

*Phys. Rev. Lett.* 71, 13:1973 (1993)

Cleo Collaboration

*hep-ex/9701008* (1997)

# References for T-odd correlation

I.Bigi *hep-ph/0107102* (2001)

G.Valencia *Phys. Rev. D* 39,11:3339 (1989)

E.Golowich, G.Valencia *Phys. Rev. D* 40,112 (1989)

B.Kayser *Nucl. Phys B (Proc. Suppl.)* 13, 487:490 (1990)

W.Bensalem, D.London *hep-ph/0005018* (2000)

J.G.Korner, K.Schilcher, Y.L.Wu *Phys. Lett. B* 242, 1:119 (1990)

B.Kayser *Nucl.Phys. B (Proc. Suppl.)* 13, 487 (1990)